

## THE DERIVATIVE OF A CONTINUED FRACTION

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### Abstract

The paper considers second order continued fractions associated with (I) the Psi function  $\psi(z)$ , (II) the continued fraction component in  $\ln \Gamma(z)$  due to Stieltjes. The second order sequences  $k_s^*/k_s$  provide approximants, some of which are remarkably close. In addition a series form for the convergent  $\chi_s/\omega_s$  associated with a continued fraction provides an expression for the derivatives of a continued fraction. The implementation uses a Maple code for derivatives.

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# 1 Introduction

We have discussed this subject in our paper Shenton and Bowman (2005, p.21). For a continued fraction it is assumed that

$$\frac{b_1}{z+} \frac{b_2}{1+} \frac{b_3}{z+} \frac{b_4}{1+} \dots = \int_0^\infty \frac{d\sigma(u)}{u+z} \quad (\Re(z) > 0)$$

the integral of Stieltjes form, the partial numerators being real and positive. The notation for the continued fraction is

$$\frac{b_1}{z + \frac{b_2}{z + \frac{b_3}{z + \frac{b_4}{z + \dots}}}}$$

In general the partial numerators are positive reals.

In this paper we introduce new examples of second order continued fractions. We are given a series

$$\frac{c_0}{z} - \frac{c_1}{z^2} + \frac{c_3}{z^2} - \frac{c_4}{z^4} + \dots$$

which in some domain may be convergent or divergent. A good reference is given in Borel (1928); see also Bowman and Shenton (1989).

To high light the basic procedures we give a simple example. For the Psi function  $\psi(\cdot)$  we have the continued fraction expression

$$\psi(z) = \ln z - \frac{1}{2z} - \frac{a_0}{z^2+} \frac{a_1}{1+} \frac{a_2}{z^2+} \dots \quad (\Re(z) > 0) \quad (1)$$

where

$$a_0 = \frac{1}{12}, \quad a_1 = \frac{1}{10}, \quad a_2 = \frac{79}{210}, \quad a_3 = \frac{1205}{1659}, \quad a_4 = \frac{262445}{209429}, \quad a_5 = \frac{2643428417511}{1429053441530}.$$

There are other interesting forms of (1).

(a) Integral form

$$\psi(z) = \ln z - \frac{1}{2z} - \int_0^\infty \frac{du}{(z^2 + u)(e^{2\pi\sqrt{u}} - 1)}$$

given by Shenton and Bowman (1971, p.552). The integral is Stieltjes and subscribes to the form

$$\int_0^\infty \frac{d\sigma(u)}{u+z^2}$$

which arises for a class of continued fraction (Wall, 1948). It suggests a continued fraction

$$\frac{d_1}{z^2+} \frac{d_2}{1+} \frac{d_3}{z^2+} \frac{d_4}{1+} \dots$$

(b) Asymptotic series

$$\psi(z) = \ln z - \frac{1}{2z} - \sum_{n=1}^{\infty} \frac{B_{2n}}{2nz^{2n}} \quad (z \rightarrow \infty \text{ in } [\arg] < \pi)$$

and in more detail

$$\psi(z) \sim \ln z - \frac{1}{2z} - \frac{1}{12z^2} + \frac{1}{20z^4} - \frac{1}{252z^6} + \frac{1}{210z^8} - \frac{1}{12z^{10}} \dots,$$

this series, extended, may be used to set up the continued fraction using Wall's algorithm (Wall, 1948, p.194).

We now return to expression (1) for  $\psi(z)$  and differentiate with respect to  $z$ , finding for the derivative of  $\psi(z)$

$$\psi_1(z) = \frac{1}{z} + \frac{1}{2z^2} - \frac{d}{dz} \left( \frac{a_0}{z^2+} \frac{a_1}{1+} \frac{a_2}{z^2+} \frac{a_3}{1+} \dots \right).$$

Thus

$$\psi_1(z) - \frac{1}{z} - \frac{1}{2z^2} = -\frac{d}{dz} \left( \frac{a_0}{z^2+} \frac{a_1}{1+} \frac{a_2}{z^2+} \frac{a_3}{1+} \dots \right)$$

so that the second order continued fraction

$$\frac{d}{d\omega} \left( \frac{a_0}{\omega+} \frac{a_1}{1+} \frac{a_2}{\omega+} \frac{a_3}{1+} \dots \right)$$

has convergent  $k_s^*/k_s(\omega)$ ,  $s = 1, 2, \dots$ , to the function  $F(z)$  where

$$F(z) = \frac{1}{2z} \left\{ \psi_1(z) - \frac{1}{z} - \frac{1}{2z^2} \right\}.$$

We shall define  $F(z)$  to be the limiting value of the 2nd order continued function sequence  $k_s^*/k_s$ . The algorithm of deriving  $k_s^*/k_s$  is given in Shenton (1957), and Shenton and Bowman (2005).

For examples, if  $z = 1$ , then  $k_s^*(1)/k_s(1)$ ,  $s = 1, 2, \dots$  converges to  $\frac{1}{2} \left( \frac{\pi^2}{6} - \frac{3}{2} \right) = 0.072$ . Moreover  $k_0^* = 0$ ,  $k_1^* = k_2^* = 1/12$ .  $k_1 = 1$ ,  $k_2 = 1 + \frac{1}{10} \left( 2 + \frac{1}{10} + \frac{79}{210} \right)$ . The approximants are  $k_1^*/k_1 = 0.08333 \dots$ ,  $k_2^*/k_2 = 0.067$ .

Table 1 Approximants for  $\psi_1(1)$ , values of  $k_s^*/k_s$

| $s$         | 1       | 2       | 3       | 4       | 5       | converges to  |
|-------------|---------|---------|---------|---------|---------|---|
| $k_s^*/k_s$ | 0.08333 | 0.06679 | 0.07492 | 0.07088 | 0.07336 | $\frac{1}{2} \left( \frac{\pi^2}{6} - \frac{3}{2} \right)$ or 0.07247 |
| $\psi_1(1)$ | 1.66667 | 1.35889 | 1.64984 | 1.64175 | 1.64671 | 1.64493   |

In the table the odd convergents corresponding to 1, 3, 5, form a monotonic decreasing sequence; similarly the even convergents, corresponding to 2, 4 form a monotonic increasing sequence. Altogether the approximations are very satisfactory.

## 2 Examples

### Example 1

In Shenton and Bowman (1971) we have given continued fractions for derivatives of the Psi functions  $\psi_m(z)$ ,  $m = 0, 1, 2, \dots$ , with  $\psi(z) = d \ln \Gamma(z)/dz$ . For  $\psi_m(z)$  itself the first those partial numerators  $C_1^{(m)}$ ,  $C_2^{(m)}$ ,  $C_3^{(m)}$  we stated explicitly. Looking at the continued fractions listed, it turns out that only one, that of  $\psi_2(z)$  has all partial numerators defined. Thus the continued fraction is

$$\frac{1}{z^2+} \frac{p_1}{1+} \frac{q_1}{z^2+} \frac{p_2}{1+} \frac{q_2}{z^2+} \dots$$

where

$$p_s = \frac{s^2(s+1)}{4s+2}, \quad q_s = \frac{s(s+1)^2}{4s+2}.$$

The example is due to Stieltjes (1918, p.388) An obvious question is why  $\psi_2(z)$  and its continued fraction? ‘

$$\psi_2(z) = -\frac{1}{z^2} - \frac{1}{z^3} - \frac{1}{2z^2} \left\{ \frac{1}{z^2+} \frac{p_1}{1+} \frac{q_1}{z^2+} \frac{p_2}{1+} \frac{q_2}{z^2+} \dots \right\} \quad (\Re(z) > 0) \quad (2)$$

(Note: There is a typo error in Shenton and Bowman (1971, p.548, (7b)); the first term on the right should be negative).

From (2),

$$\left( \psi_2(z) + \frac{1}{z^2} + \frac{1}{z^3} \right) 2z^2 = -\frac{1}{z^2+} \frac{p_1}{1+} \frac{q_1}{z^2+} \dots$$

differentiate with respect to  $z$

$$2 \frac{d}{dz} \left( z^2 \psi_2(z) + 1 + \frac{1}{z} \right) = 2z \frac{d}{d\omega} \left( \frac{1}{\omega+} \frac{p_1}{1+} \frac{q_1}{\omega+} \frac{p_2}{1+} \frac{q_2}{\omega+} \dots \right) \quad (\omega = z^2)$$

resulting

$$\left( 2\psi_2(z) + z\psi_3(z) - \frac{1}{z^3} \right) = \lim \frac{k_s^*(\omega)}{k_s(\omega)} \quad (\omega = z^2)$$

and

$$\frac{k_1^*}{k_1} = \frac{1}{\omega^2}, \quad \frac{k_2^*}{k_2} = \frac{1}{\omega^2 + p_1(2\omega + q_1 + p_1)},$$

the first should be an upper bound and the second is a lower bound. Examples are given in Table 2.

Table 2 Values of  $k_s^*/k_s$  of example 1

| $z$ | 1          | 2          | 3          | 4         | 5          | converges to |
|-----|------------|------------|------------|-----------|------------|--------------|
| 1/2 | 16.0000000 | 1.77777778 | 12.6137566 | 3.0634024 | 10.9613493 | 7.046952     |
| 1   | 1.0000000  | 0.5000000  | 0.7916666  | 0.7322530 | 0.6134259  | 0.685711     |
| 2   | 0.0625000  | 0.0526316  | 0.0552326  | 0.0544314 | 0.0547382  | 0.054651     |

### Example 2

$$\ln \Gamma(z) = \left(z - \frac{1}{2}\right) \ln z - z + \frac{1}{2} \ln(2\pi) + \frac{a_0}{z+} \frac{a_1}{z+} \frac{a_2}{z+} \dots \quad (\Re(z) > 0)$$

where  $a_0$  to  $a_6$ , are given in Wall (1948).

The continued fraction may be written

$$\frac{\alpha_1 z}{z^2+} \frac{\alpha_2}{1+} \frac{\alpha_3}{z^2+} \frac{\alpha_4}{1+} \frac{\alpha_5}{z^2+} \frac{\alpha_6}{1+} \dots (\Re(z) > 0),$$

so we write

$$\frac{1}{z} \left\{ \ln \Gamma(z) - \left(z - \frac{1}{2}\right) \ln z + z - \frac{\ln(2\pi)}{2} \right\} = \frac{\alpha_1}{\omega+} \frac{\alpha_2}{1+} \frac{\alpha_3}{\omega+} \frac{\alpha_4}{1+} \dots \quad (\omega = z^2)$$

and the right hand side is in the correct form to agree with the 2nd order continued fraction. Differentiate with respect to  $z$ , we have

$$\frac{d}{dz} \frac{1}{z} \left\{ \ln \Gamma(z) - \left(z - \frac{1}{2}\right) \ln z + z - \frac{\ln(2\pi)}{2} \right\} = 2z \frac{d}{d\omega} \left\{ \frac{\alpha_1}{\omega+} \frac{\alpha_2}{1+} \frac{\alpha_3}{\omega+} \frac{\alpha_4}{1+} \dots \right\}.$$

Allowing for a change in sign when we consider derivatives, the sequences  $k_s^*/k_s$  are approximants to

$$\begin{aligned} & -\frac{1}{2z} \frac{d}{dz} \left\{ \frac{\ln \Gamma(z)}{z} - \left(1 - \frac{1}{2z}\right) \ln z + 1 - \frac{\ln(2\pi)}{2z} \right\} \\ & = -\frac{1}{2z} \left\{ \frac{\psi_1(z)}{z} - \frac{\ln \Gamma(z)}{z^2} - \frac{\ln(z)}{2z^2} - \left(\frac{1}{z} - \frac{1}{2z^2}\right) + \frac{\ln(2\pi)}{2z^2} \right\} \end{aligned}$$

with

$$k_{1*} = 1/12 = k_2^*, \quad k_1 = \omega^2, \quad k_2 = \omega^2 + \alpha_2(2\omega + \alpha_3 + \alpha_2).$$

and when  $z = 1/2$  becomes

$$-\left\{2\psi_1\left(\frac{1}{2}\right) + 4 \ln 2\right\} = 2\gamma$$

Take  $z = 1$ . The 2nd order continued fraction ( $k_s^*/k_s$ ) converges to the function

$$F(1) = -\frac{1}{2} \left\{ \psi(1) - \left(1 - \frac{1}{2}\right) - \frac{\ln(2\pi)}{2} \right\} = 0.07908$$

Table 3 Values of  $k_s^*/k_s$  of example 2

| $z$ | 1         | 2         | 3         | 4         | 5         | converges to |
|-----|-----------|-----------|-----------|-----------|-----------|--------------|
| 1/2 | 1.3333333 | 0.9395973 | 1.2501182 | 1.0464910 | 1.1833200 | $2\gamma$    |
| 1   | 0.0833333 | 0.0774363 | 0.0797778 | 0.0787481 | 0.0792193 | 0.07908      |
| 2   | 0.0052083 | 0.0051200 | 0.0051306 | 0.0051288 | 0.0051292 | 0.0051291    |

### 3 A determinantal reduction formula for continued fractions

Consider the continued fraction  $\frac{c_0}{z+} \frac{c_1}{z+} \frac{c_2}{z+} \dots$  with convergents  $\frac{\chi_s}{\omega_s}$ ,  $s = 1, 2, \dots$ . Then, for example,

$$\begin{aligned} \chi_1 &= c_0, & \chi_2 &= zc_1, \\ \omega_1 &= z, & \omega_2 &= z + c_1. \quad (\chi_0 = 0, \omega_0 = 1). \end{aligned}$$

We may therefore consider the determinant

$$\begin{vmatrix} \chi_s & \chi_{s+1} \\ \omega_s & \omega_{s+1} \end{vmatrix}$$

using the recurrence formulas. Expanding the determinant leads to

$$\frac{\chi_s}{\omega_s} = \frac{\chi_0}{\omega_0} + \left(\frac{\chi_1}{\omega_1} - \frac{\chi_0}{\omega_0}\right) + \dots + \left(\frac{\chi_s}{\omega_s} - \frac{\chi_{s-1}}{\omega_{s-1}}\right) + \dots +$$

at least formally. Hence

$$\frac{\chi_s}{\omega_s} = \frac{c_0}{\omega_0\omega_1} + \frac{c_0c_1}{\omega_1\omega_2} + \frac{c_0c_1c_2}{\omega_2\omega_3}. \quad (3)$$

#### Example 3

The Laplace continued fraction is

$$\frac{1}{z+} \frac{1}{z+} \frac{2}{z+} \frac{3}{z+} \dots \quad (z > 0)$$

so  $\omega_1 = z$ ,  $\omega_2 = z^2 + z$ ,  $\omega_3 = z^3 + 3z$ ,  $\omega_4 = z^4 + 6z^2 + 3$ . A simple formula is  $\omega_s = e^{-\frac{1}{2}d_x^2} x^2$  the polynomial being Hermite.

Now using (3) we may set up an expression for derivatives of a continued fraction using the Maple symbolic code.

## 4 Conclusion

We now pay attention to the function  $F(z)$  which represents the value to which the 2nd order continued fractions converges. If  $z$  is real and positive is  $F(z)$  positive? For the first example we knew (Shenton and Bowman, 1971)

$$\begin{aligned} F(z) &= \frac{1}{2z} \left\{ \psi_1(z) - \frac{1}{z} - \frac{1}{2z^2} \right\} \\ &= \frac{1}{2z} \frac{2\pi}{3} \int_0^\infty \frac{y\sqrt{x}dx}{xz^2(y-1)^2} \quad (y = e^{2\pi\sqrt{x}}) \end{aligned}$$

which is positive for  $z > 0$ .

Now consider the expression for  $F(z)$  arising from the continued fraction from  $\ln \Gamma(z)$ . In the sequel for this case

$$\begin{aligned} F(z) &= -\frac{1}{2z} \frac{d}{dz} \left\{ \frac{\ln \Gamma(z)}{z} - \left(1 - \frac{1}{2z}\right) \ln z + 1 - \frac{\ln(2\pi)}{2z} \right\} \\ &= -\frac{1}{2z} \left\{ \frac{\psi(z)}{z} - \frac{\ln \Gamma(z)}{z^2} - \frac{\ln z}{2z^2} - \left(\frac{1}{z} - \frac{1}{2z^2}\right) - \frac{\ln(2\pi)}{2z^2} \right\} \end{aligned}$$

We have been unable to prove this is positive, for  $z$  real and positive. So we test out several cases.

Case with  $z = \frac{1}{2}$ :

$$F(z) = \{-2\psi(2) - 2 \ln \pi - 2 \ln \frac{1}{2} + 2 \ln(2\pi)\} = 2\gamma \quad (\gamma = \text{Euler's constant}).$$

The upper bound is 4/3. It is quite remarkable that in the Handbook of Mathematical Functions.... section 6.3.3, we have the single entry (also see "dlmf.nist.gov", 5.4.13)

$$\psi\left(\frac{1}{2}\right) = -\gamma - 2 \ln 2 = 1.9635100260\dots$$

As far as we can tell it is not referenced. What a remarkable piece of luck?

Now take  $F(1)$ : We have

$$F(1) = -\frac{1}{2} \left\{ -\gamma - \frac{1}{2} + \frac{\ln(2\pi)}{2} \right\} = 0.0058.$$

The upper bound is  $1/12$ .

For  $z = 2$ ,

$$\begin{aligned} F(2) &= -\frac{1}{4} \left\{ -\frac{\psi(2)}{2} - \frac{\ln 2}{8} - \frac{3}{8} + \frac{\ln(2\pi)}{8} \right\} \\ &= -\frac{1}{4} \left\{ \frac{1}{2}(1 - \gamma) - \frac{3}{8} + \frac{\ln(2\pi)}{8} \right\} = 0.0051\dots \end{aligned}$$

The upper bound is 0.0208.

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