



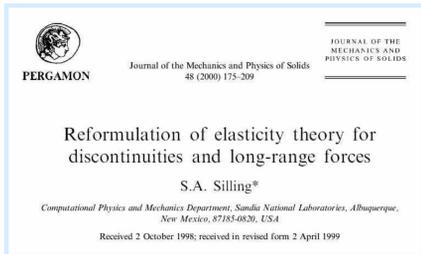
Mathematical and numerical analysis of peridynamics for multiscale materials modeling

Qiang Du (PSU), Max Gunzburger (FSU), Rich Lehoucq (SNL)



INTRODUCTION

- Silling proposes the peridynamic nonlocal continuum theory



Reformulation of elasticity theory for discontinuities and long-range forces

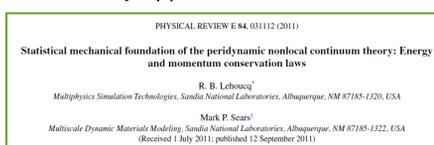
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- ✓ **Motivation:** materials undergoing discontinuous deformation
- ✓ **Goal:** Dynamic material failure simulations
- ✓ **Key:** Nonlocal model of force via integral operators
- ✓ Generalization to "State-based" theory proposed in 2007

- A statistical mechanical basis for nonlocality appears in



Statistical mechanical foundation of the peridynamic nonlocal continuum theory: Energy and momentum conservation laws

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- ✓ Important conclusion is that nonlocality is intrinsic to continuum balance laws

- Recent review



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Peridynamic Theory of Solid Mechanics

S.A. Silling* and R.B. Lehoucq†

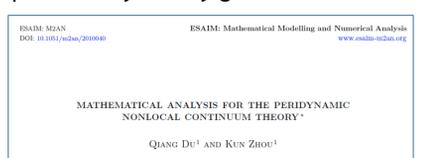
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- ✓ Balance of energy; second law of thermodynamics

OBJECTIVES

- Mathematical analysis for the peridynamic continuum theory; preliminary theory given in



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MATHEMATICAL ANALYSIS FOR THE PERIDYNAMIC NONLOCAL CONTINUUM THEORY*

QIANG DU¹ AND KUN ZHOU²



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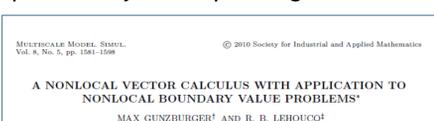
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MATHEMATICAL AND NUMERICAL ANALYSIS OF LINEAR PERIDYNAMIC MODELS WITH NONLOCAL BOUNDARY CONDITIONS*

KUN ZHOU¹ AND QIANG DU²

- ✓ In particular, provide a mathematical analysis for the state-based theory

- Describe *volume-constraints*, the nonlocal analogue of boundary conditions
- Analysis facilitated by the development of a *nonlocal vector calculus*; preliminary development given in



MULTISCALE MODEL. SIMUL.

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A NONLOCAL VECTOR CALCULUS WITH APPLICATION TO NONLOCAL BOUNDARY VALUE PROBLEMS*

MAX GUNZBURGER¹ AND R. B. LEHOUCQ²

- Short-term goal: linear peridynamic materials, diffusion, numerical analysis
- Long-term goal: nonlinear peridynamic materials, coupling with the classic theory and molecular dynamics

OVERVIEW

- Balance of linear momentum

$$\rho(x, t) \dot{y}(x, t) = \int_{\mathbb{R}^3} (\mathbf{t}(x', x, t) - \mathbf{t}(x, x', t)) dx' + b(x, t)$$

- Balance of linear momentum is equivalent to or action-reaction.

$$\int_{\Omega_1} \int_{\Omega_2} (\mathbf{t}(x', x, t) - \mathbf{t}(x, x', t)) dx' dx + \int_{\Omega_2} \int_{\Omega_1} (\mathbf{t}(x', x, t) - \mathbf{t}(x, x', t)) dx' dx = 0$$

$$\forall \Omega_1 \& \Omega_2, \Omega_1 \cap \Omega_2 = \emptyset$$

- Nonlocal because force may be nonzero even when Ω_1 and Ω_2 are not in contact

$$\underbrace{\int_{\Omega} \rho(x, t) \dot{y}(x, t) dx}_{\text{Rate of change of momentum}} = \underbrace{\int_{\Omega} \int_{\mathbb{R}^3 \setminus \Omega} (\mathbf{t}(x', x, t) - \mathbf{t}(x, x', t)) dx' dx}_{\text{Force exerted upon } \Omega} + \underbrace{\int_{\Omega} b(x, t) dx}_{\text{External force upon } \Omega}$$

- Balance of energy

$$\underbrace{\frac{d}{dt} \int_{\Omega} \varepsilon(x, t) dx}_{\text{Internal energy}} + \underbrace{\frac{1}{2} \frac{d}{dt} \int_{\Omega} \dot{y}(x, t) \cdot \dot{y}(x, t) dx}_{\text{Kinetic energy over } \Omega} = \underbrace{\int_{\Omega} \int_{\mathbb{R}^3 \setminus \Omega} (\mathbf{t}(x', x) \cdot \dot{y}(x', t) - \mathbf{t}(x, x') \cdot \dot{y}(x, t)) dx' dx}_{\text{Supplied power to } \Omega} + \underbrace{\int_{\Omega} \int_{\mathbb{R}^3 \setminus \Omega} q(x', x) dx' dx}_{\text{Supplied thermal power to } \Omega}$$

$$\dot{\varepsilon}(x, t) = \underbrace{\int_{\mathbb{R}^3} \mathbf{t}(x', x) \cdot (\dot{y}(x', t) - \dot{y}(x, t)) dx'}_{\text{Absorbed power density}} + \underbrace{\int_{\mathbb{R}^3} q(x', x) dx'}_{\text{Thermal power}}$$

- Constitutive relations

Define the deformation state $\underline{Y}[x, t] \langle x' - x \rangle := y(x', t) - y(x, t) \quad \forall x'$ so that the *Force*

State $\underline{T}[x, t] = \hat{\underline{T}}(\underline{Y}[x, t])$ depends collective motion

$$\mathbf{t}(x', x, t) = \underline{T}[x, t] \langle x' - x \rangle = \hat{\underline{T}}(\underline{Y}[x, t]) \langle x' - x \rangle$$

The needed relations can be written as

$$q(x', x, t) = \underline{Q}[x, t] \langle x' - x \rangle = \hat{\underline{Q}}(\underline{Y}[x, t]) \langle x' - x \rangle$$

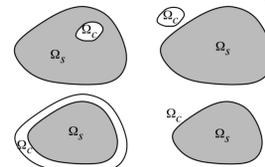
- Can now discuss the well-posedness of the balance laws; in particular for linear materials
- Can also show that the second law of thermodynamic is not violated for classes of materials

CASE STUDY: NONLOCAL DIFFUSION

- Special case of the balance of energy of linear diffusion (modeling anomalous diffusion) serves to introduce the nonlocal vector calculus and provide well-posedness
- Consider the nonlocal Dirichlet problem (steady-state nonlocal diffusion)

$$\mathcal{D}(\Theta \mathcal{D}^*) u = f \quad \text{on } \Omega_s$$

$$u = g \quad \text{on } \Omega_c$$



- The solution constrained over the volume Ω_c
- A volume-constrained problem is the nonlocal analogue of a boundary value problem

$$(\mathcal{D}f)(x) := \int_{\mathbb{R}^n} (f(x, y) + f(y, x)) \cdot \alpha(x, y) dy, \quad \alpha(x, y) = -\alpha(y, x)$$

$$(\mathcal{D}^* u)(x, y) := -(u(y, t) - u(x, t)) \alpha(x, y)$$

$$\mathcal{D}(\Theta \mathcal{D}^* u)(x) = 2 \int_{\mathbb{R}^n} (u(y, t) - u(x, t)) \alpha(x, y) \cdot \Theta(x, y) \alpha(x, y) dy$$

- \mathcal{D} and \mathcal{D}^* are the nonlocal divergence and it's adjoint; the operator $\mathcal{D}\mathcal{D}^*$ is the nonlocal Laplacian
- The kernel $\alpha\Theta\alpha$ determines the regularity for the volume-constrained problem; an integrable kernel implies that $\mathcal{D}\Theta\mathcal{D}^* u : L^2(\Omega_s \cup \Omega_c) \rightarrow L^2(\Omega_s \cup \Omega_c)$; no smoothing of the data
- Fractional smoothing occurs for $\alpha(x, y)\Theta(x, y)\alpha(x, y) \sim |x-y|^{-n-2s}, 0 < s < 1$ because then $\mathcal{D}\Theta\mathcal{D}^* u : H^s(\Omega_s \cup \Omega_c) \rightarrow H^s(\Omega_s \cup \Omega_c)$
- Can extend to nonlocal Neumann, Robin problems; also consider peridynamic linear elastic volume-constrained problems

SUMMARY

- Peridynamic continuum theory reformulated using a nonlocal vector calculus
- Variational formulation leads to the well-posedness of the
 - ✓ peridynamic equilibrium equation for linear isotropic solids
 - ✓ nonlocal linear diffusion
- Deformation can be discontinuous
- Volume-constraints, the nonlocal analogue of boundary conditions, are crucial

Related work

- Probabilistic interpretation for nonlocal linear diffusion
- Nonlocal, nonlinear advection
- Finite element formulation; see



- Two Ph.d theses; Pablo Seleson (FSU) and Nate Burch (CSU)

PUBLICATIONS

- A nonlocal vector calculus, nonlocal volume-constrained problems, and nonlocal balance laws* SAND 2010-8353J (Q. Du, M. Gunzburger, R. Lehoucq, K. Zhou).
- An approach to nonlocal, nonlinear advection* SAND 2011-3164J (Q. Du, J. Kamm, R. Lehoucq, M. Parks).
- Analysis and approximation of nonlocal diffusion problems with volume constraints* SAND 2011-3168J (Q. Du, M. Gunzburger, R. Lehoucq, K. Zhou)
- Application of a nonlocal vector calculus to the analysis of linear peridynamic materials* SAND 2011-3870J (Q. Du, M. Gunzburger, R. Lehoucq, K. Zhou)
- A posteriori error analysis of finite element method for linear nonlocal diffusion and peridynamic models* (Q. Du, L. Ju, L. Tian, K. Zhou)

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